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Abstract. We deal with initial-boundary value problems describing vertical vibrations of viscoelastic von Kármán-Donnell shells with a rigid inner obstacle. The short memory (Kelvin-Voigt) material is considered. A weak formulation of the problem is in the form of the hyperbolic variational inequality. We solve the problem using the penalization method.

Keywords: Von Kármán-Donnell shell, unilateral dynamic contact, viscoelasticity, solvability, penalty approximation

1   Introduction

Contact problems represent an important but complex topic of applied mathematics. Its complexity profounds if the dynamic character of the problem is respected. For elastic problems there is only a very limited amount of results available (cf. [3] and there cited literature). Viscosity makes possible to prove the existence of solutions for a broader set of problems for membranes, bodies as well as for linear models of plates. The presented results extend the research made in [2], where the problem for a viscoelastic short memory von Kármán plate in a dynamic contact with a rigid obstacle was considered. Our results also extend the research made for the quasistatic contact problems for viscoelastic shells (cf. [1]). A thin isotropic shallow shell occupies the domain

$$G = \{(x,z) \in R^3 : \ x = (x_1, x_2) \in \Omega, \ |z - \mathcal{Z}| < h/2\},$$

where $h > 0$ is the thickness of the shell, $\Omega \subset \mathbb{R}^2$ is a bounded simply connected domain in $\mathbb{R}$ with a sufficiently smooth boundary Γ. We set $I \equiv (0, T)$ a bounded time interval, $Q = I \times \Omega, \ S = I \times \Gamma$. The unit outer normal vector is denoted by $\mathbf{n} = (n_1, n_2), \ \tau = (-n_2, n_1)$ is the unit tangent vector. The displacement is denoted by $\mathbf{u} \equiv (u_i)$. The strain tensor is defined as

$$\varepsilon_{ij}(\mathbf{u}) = \frac{1}{2}(\partial_i u_j + \partial_j u_i + \partial_i u_3 \partial_j u_3) - k_{ij} u_3 - x_3 \partial_{ij} u_3, \ i, j = 1, 2$$

with $k_{12} = k_{21} = 0$ and the curvatures $k_{ii} > 0$, $i = 1, 2$.

Further, we set

$$[u, v] \equiv \partial_{11} u \partial_{22} v + \partial_{22} u \partial_{11} v - 2 \partial_{12} u \partial_{12} v.$$

In the sequel, we denote by $W_p^k(M)$, $k \geq 0$, $p \in [1, \infty]$ the Sobolev spaces defined on a domain or an appropriate manifold M. By $\mathring{W}_p^k(M)$ the spaces with zero traces are denoted. If $p = 2$ we use the notation $H^k(M)$, $\mathring{H}^k(M)$. The duals to $\mathring{H}^k(M)$ are denoted by $H^{-k}(M)$. For the anisotropic spaces $W_p^k(M)$, $k = (k_1, k_2) \in \mathbb{R}_+^2$, k_1 is related with the time variable while k_2 with the space variables. We shall use also the Bochner-type spaces $W_p^k(I; X)$ for a time interval I and a Banach space X. Let us remark that for $k \in (0, 1)$ their norm is defined by the relation

$$\|w\|_{W_p^k(I;X)}^p \equiv \int_I \|w(t)\|_X^p \, dt + \int_I \int_I \frac{\|w(t) - w(s)\|_X^p}{|s - t|^{1+kp}} \, ds \, dt.$$

By $C(M)$ we denote the spaces of continuous functions on a (possibly relatively) compact manifold M. They are equipped with the max-norm. Analogously the spaces $C(M; X)$, are introduced for a Banach space X. The following generalization of the Aubin's compactness lemma verified in [4] Theorem 3.1 will be essentially used:

Lemma 1 *Let $B_0 \hookrightarrow\hookrightarrow B \hookrightarrow B_1$ be Banach spaces, the first reflexive and separable. Let $1 < p < \infty$, $1 \leq r < \infty$. Then*

$$W \equiv \{v; \ v \in L_p(I; B_0), \ \dot{v} \in L_r(I, B_1)\} \hookrightarrow\hookrightarrow L_p(I; B).$$

2   Short memory material

2.1   Problem formulation

Employing the Einstein summation, the constitutional law has the form

$$\sigma_{ij}(\mathbf{u}) = \frac{E_1}{1 - \mu^2} \partial_t \big((1-\mu)\varepsilon_{ij}(\mathbf{u}) + \mu \delta_{ij} \varepsilon_{kk}(\mathbf{u})\big) + \frac{E_0}{1 - \mu^2} \big((1-\mu)\varepsilon_{ij}(\mathbf{u}) + \mu \delta_{ij} \varepsilon_{kk}(\mathbf{u})\big).$$

The constants E_0, $E_1 > 0$ are the Young modulus of elasticity and the modulus of viscosity, respectively. We shall use the abbreviation $b = h^2/(12\varrho(1 - \mu^2))$, where $h > 0$ is the shell thickness and ϱ is the density of the material. We involve the rotation inertia expressed by the term $a\Delta\ddot{u}$ in the first equation of the considered system with $a = \frac{h^2}{12}$. It will play the crucial role in the deriving a strong convergence of the sequence of velocities $\{\dot{u}_m\}$ in the appropriate space. We assume the shell clamped on the boundary. We generalize the dynamic elastic model due to the von Kármán-Donnell theory mentioned in [6]. The classical formulation for the deflection $u_3 \equiv u$ and the Airy stress function v is then the

initial-value problem

$$
\left.\begin{aligned}
&\ddot{u} + a\Delta\ddot{u} + b(E_1\Delta^2\dot{u} + E_0\Delta^2 u) - [u,v] - \Delta_k * v = f + g,\\
&u - \Psi \geq 0,\ g \geq 0,\ (u - \Psi)g = 0,\\
&\Delta^2 v + E_1\partial_t(\tfrac{1}{2}[u,u] + k_{11}\partial_{22}u + k_{22}\partial_{11}u)\\
&+E_0(\tfrac{1}{2}[u,u] + \Delta_k u) = 0
\end{aligned}\right\} \text{ on } Q, \quad (1)
$$

$$
u = \partial_n u = v = \partial_n v = 0 \text{ on } S, \tag{2}
$$

$$
u(0,\cdot) = u_0,\ \dot{u}(0,\cdot) = u_1 \text{ on } \Omega. \tag{3}
$$

The obstacle function $\Psi \in L_\infty(\Omega)$ is fulfilling $0 < U_0 \leq u_0 - \Psi$ in Ω and

$$
\Delta_k u \equiv \partial_{11}(k_{22}u) + \partial_{22}(k_{11}u), \tag{4}
$$

$$
\Delta_k^* v \equiv k_{22}\partial_{11}v + k_{11}\partial_{22}v. \tag{5}
$$

We define the operators $L : H^2(\Omega) \to \mathring{H}^2(\Omega)$, $\Phi : H^2(\Omega) \times H^2(\Omega) \to \mathring{H}^2(\Omega)$ by uniquely solved equations

$$
(\Delta Lu, \Delta w) \equiv (\Delta_k u, w)\, \forall w \in \mathring{H}^2(\Omega), \tag{6}
$$

$$
(\Delta\Phi(u,v), \Delta w) \equiv ([u,v], w)\, \forall w \in \mathring{H}^2(\Omega). \tag{7}
$$

with the inner product $(\cdot, \cdot)$ in the space $L_2(\Omega)$. The operator L is linear and compact. The bilinear operator Φ is symmetric and compact. Moreover due to Lemma 1 from [5] $\Phi : H^2(\Omega)^2 \to W_p^2(\Omega)$, $2 < p < \infty$ and

$$
\|\Phi(u,v)\|_{W_p^2(\Omega)} \leq c\|u\|_{H^2(\Omega)}\|v\|_{W_p^1(\Omega)}\ \forall u \in H^2(\Omega),\ v \in W_p^1(\Omega). \tag{8}
$$

We have also $L : H^2(\Omega) \mapsto W_p^2(\Omega)$, $2 < p < \infty$ and

$$
\|Lu\|_{W_p^2(\Omega)} \leq c\|u\|_{H^2(\Omega)}\ \forall u \in H^2(\Omega). \tag{9}
$$

For $u, y \in L_2(I; H^2(\Omega))$ we define the bilinear form A by

$$
A(u,y) := b\big(\partial_{kk}u\partial_{kk}y + \mu(\partial_{11}u\partial_{22}y + \partial_{22}u\partial_{11}y) + 2(1-\mu)\partial_{12}u\partial_{12}y\big).
$$

We introduce shifted cone $\mathcal{K}$ by

$$
\mathcal{K} := \{y \in H^{1,2}(Q);\ \dot{y} \in L_2(I, \mathring{H}^1(\Omega);\ y \geq \Psi\}. \tag{10}
$$

Then the variational formulation of the problem (1-3) has the form of

Problem $\mathcal{P}$. *Find $u \in \mathcal{K}$ such that $\dot{u} \in L_2(I; \mathring{H}^2(\Omega))$ and*

$$
\begin{aligned}
&\int_Q (E_1 A(\dot{u}, y - u) + E_0 A(u, y - u))\, dx\, dt\\
&+ \int_Q [u, E_1\partial_t(\tfrac{1}{2}\Phi(u,u) + Lu) + E_0(\tfrac{1}{2}\Phi(u,u) + Lu)](y - u)\, dx\, dt\\
&+ \int_Q \Delta_k\big(E_1\partial_t(\tfrac{1}{2}\Phi(u,u) + Lu) + E_0(\tfrac{1}{2}\Phi(u,u) + Lu)\big)(y - u)\, dx\, dt\\
&- \int_Q (a\nabla\dot{u}\cdot\nabla(\ddot{y} - \dot{u}) + \dot{u}(\dot{y} - \dot{u}))\, dx\, dt\\
&+ \int_\Omega (a\nabla\dot{u}\cdot\nabla(y - u) + \dot{u}(y - u))(T, \cdot)\, dx\\
&\geq \int_\Omega (a\nabla u_1\cdot\nabla(y(0,\cdot) - u_0) + u_1(y(0,\cdot) - u_0))\, dx\\
&+ \int_Q f(y_1 - u)\, dx\, dt\ \forall y \in \mathcal{K}.
\end{aligned} \tag{11}
$$

2.2   The penalization

For any $\eta > 0$ we define the *penalized problem*

Problem $\mathcal{P}_\eta$. *Find $u \in H^{1,2}(Q)$ such that $\dot{u} \in L_2(I; \mathring{H}^2(\Omega))$, $\ddot{u} \in L_2(I; \mathring{H}^1(\Omega))$,*

$$
\begin{aligned}
&\int_Q \left(\ddot{u}z + a\nabla\ddot{u} \cdot \nabla z + E_1 A(\dot{u}, z) + E_0 A(u, z) \right) dx\, dt \\
&+ \int_Q [u, E_1 \partial_t(\tfrac{1}{2}\Phi(u,u) + Lu) + E_0(\tfrac{1}{2}\Phi(u,u) + Lu)]z\, dx\, dt \\
&+ \int_Q \Delta_k \left(E_1 \partial_t(\tfrac{1}{2}\Phi(u,u) + Lu) + E_0(\tfrac{1}{2}\Phi(u,u) + Lu) \right) z\, dx\, dt \\
&= \int_Q (f + \eta^{-1}(u - \Psi)^-)z\, dx\, dt \ \forall z \in L_2(I; H^2(\Omega))
\end{aligned}
\tag{12}
$$

and the conditions (3) remain valid.

Lemma 2 *Let $f \in L_2(Q)$, $u_0 \in \mathring{H}^2(\Omega)$, and $u_1 \in \mathring{H}^1(\Omega)$. Then there exists a solution u of the problem $\mathcal{P}_\eta$.*

Proof. Let us denote by $\{w_i \in \mathring{H}^2(\Omega); i = 1, 2, ...\}$ a basis of $\mathring{H}^2(\Omega)$ orthonormal in $H^1(\Omega)$ with respect to the inner product

$$
(u, v)_a = \int_\Omega (uv + a\nabla u \cdot \nabla v)\, dx, \ u, v \in H^1(\Omega).
$$

We construct the Galerkin approximation u_m of a solution in a form

$$
u_m(t) = \sum_{i=1}^m \alpha_i(t)w_i, \ \alpha_i(t) \in \mathbb{R}, \ i = 1, ..., m, \ m \in N,
\tag{13}
$$

$$
\begin{aligned}
&(\ddot{u}_m(t), w_i)_a + \int_\Omega \left(E_1 A(\dot{u}_m(t), w_i) + E_0 A(u_m(t), w_i) \right) dx + \\
&\int_\Omega \Delta\left(E_1 \partial_t(\tfrac{1}{2}\Phi(u_m, u_m) + Lu_m) + E_0(\tfrac{1}{2}\Phi(u_m, u_m) + Lu_m) \right) \\
&\times \Delta(\Phi(u_m, w_i) + Lw_i)\, dx \\
&= \int_\Omega \left(f(t) + \eta^{-1}(u_m(t) - \Psi)^- \right)w_i\, dx, \ i = 1, ..., m,
\end{aligned}
\tag{14}
$$

$$
u_m(0) = u_{0m}, \ \dot{u}_m(0) = u_{1m}, \ u_{0m} \to u_0 \text{ in } \mathring{H}^2(\Omega), \ u_{1m} \to u_1 \text{ in } \mathring{H}^1(\Omega).
\tag{15}
$$

After multiplying the equation (14) by $\dot{\alpha}_i(t)$, summing up with respect to i, taking in mind the definitions of the operators Φ, L and integrating we obtain the *a priori* estimates not depending on m:

$$
\begin{aligned}
&\|\dot{u}_m\|^2_{L_2(I;\mathring{H}^2(\Omega))} + \|\dot{u}_m\|^2_{L_\infty(I;\mathring{H}^1(\Omega))} + \|u_m\|^2_{L_\infty(I;\mathring{H}^2(\Omega))} \\
&+ \|\partial_t \Phi(u_m, u_m)\|^2_{L_2(I;\mathring{H}^2(\Omega))} + \|\partial_t Lu_m\|^2_{L_2(I;\mathring{H}^2(\Omega))} \\
&+ \eta^{-1}\|(u_m - \Psi)^-\|_{L_\infty(I;L_2(\Omega))} \le c \equiv c(f, u_0, u_1).
\end{aligned}
\tag{16}
$$

Moreover the estimates (8), (9) imply

$$
\|\partial_t \Phi(u_m, u_m)\|_{L_2(I;W_p^2(\Omega))} + \|\partial_t Lu_m\|_{L_2(I;W_p^2(\Omega))} \le c_p \ \forall p > 2.
\tag{17}
$$

After multiplying the equation (14) by $\ddot{\alpha}_i(t)$, summing up and integrating we obtain the estimate of $\ddot{u}_m$

$$
\|\ddot{u}_m\|_{L_2(I;H^1(\Omega))} \le c_\eta, \ m \in \mathbb{N}.
\tag{18}
$$

Applying the estimates (16)-(18), the compact imbedding theorem and the interpolation, we obtain for any $p \in [1, \infty)$, a subsequence of $\{u_m\}$ (denoted again by $\{u_m\}$), a function u and the convergences

$$
\begin{aligned}
&\ddot{u}_m \rightharpoonup \ddot{u} \text{ in } L_2(I; H^1(\Omega)), \\
&\dot{u}_m \rightharpoonup^* \dot{u} \text{ in } L_\infty(I; \mathring{H}^1(\Omega)), \\
&\dot{u}_m \rightharpoonup \dot{u} \text{ in } L_2(I; \mathring{H}^2(\Omega)), \\
&\dot{u}_m \rightarrow \dot{u} \text{ in } L_p(I; \mathring{H}(\Omega)) \cap L_\infty(I; H^{2-\varepsilon}(\Omega)) \; \forall \varepsilon > 0, \\
&u_m \rightarrow u \text{ in } C(\bar{I}; W_p^1(\Omega), \\
&\partial_t\left(\tfrac{1}{2}\Phi(u_m, u_m) + Lu_m\right) \rightarrow \partial_t\left(\tfrac{1}{2}\Phi(u, u) + Lu\right) \text{ in } L_2(I; W_p^2(\Omega))
\end{aligned}
\tag{19}
$$

implying that a function u fulfils the identity (12). The initial conditions (3) follow due to (15) and the proof of the existence of a solution is complete.

2.3 Solving the original problem

We verify the existence theorem

Theorem 1 *Let* $f \in L_2(Q)$, $u_i \in \mathring{H}^2(\Omega)$, $i = 0, 1$, $0 < U_0 \leq u_0 - \Psi$. *Then there exists a solution of the Problem* $\mathcal{P}$.

Proof. We perform the limit process for $\eta \rightarrow 0$. We write u_η for the solution of the problem $\mathcal{P}_{1,\eta}$. The *a priori* estimates (16) imply the estimates

$$
\begin{aligned}
&\|\dot{u}_\eta\|^2_{L_2(I; \mathring{H}^2(\Omega))} + \|\dot{u}_\eta\|^2_{L_\infty(I; \mathring{H}^1(\Omega))} + \|u_\eta\|^2_{L_\infty(I; \mathring{H}^2(\Omega))} \\
&+\|\partial_t\Phi(u_\eta, u_\eta)\|^2_{L_2(I; W_p^2(\Omega))} + \|\partial_t Lu_\eta\|^2_{L_2(I; W_p^2(\Omega))} \\
&+\eta^{-1}\|(u_\eta - \Psi)^-\|_{L_\infty(I; L_2(\Omega))} \leq c_p, \; p > 2.
\end{aligned}
\tag{20}
$$

To get the crucial estimate for the penalty, we put $z = u_0 - u_\eta(t, \cdot)$ in (12) and obtain the estimate

$$
\begin{aligned}
0 \leq U_0 \int_Q \eta^{-1}(u_\eta - \Psi)^- dx\, dt &\leq \int_Q \|\eta^{-1}(u_\eta - \Psi)^-(u_0 - \Psi) dx\, dt \\
&\leq \int_Q \|\eta^{-1}(u_\eta - \Psi)^-(u_0 - u_\eta) dx\, dt \\
&= \int_Q \left(\dot{u}_\eta^2 + a|\nabla\dot{u}_\eta|^2 + A((E_1\partial_t u_\eta + E_0 u_\eta), u_0 - u_\eta)\right. \\
&+E_1\partial_t(\Delta(Lu_\eta + \tfrac{1}{2}\Phi(u_\eta, u_\eta)))\Delta(L(u_0 - u_\eta) + \Phi(u_\eta, u_0 - u_\eta)) \\
&+E_0\Delta(Lu_\eta + \tfrac{1}{2}\Phi(u_\eta, u_\eta))\Delta(L(u_0 - u_\eta) + \Phi(u_\eta, u_0 - u_\eta))\Big) dx\, dt \\
&- \int_Q f(u_0 - u_\eta)\, dx\, dt + \int_\Omega \left((\dot{u}_\eta(u_0 - u_\eta) + a\nabla\dot{u}_\eta \cdot \nabla(u_0 - u_\eta))(T, \cdot)\right) dx.
\end{aligned}
$$

Applying the *a priori* estimates (20) we obtain

$$
\|\eta^{-1}u_\eta^-\|_{L_1(Q)} \leq c(f, u_0, u_1, \Psi).
\tag{21}
$$

With respect to Dirichlet conditions we obtain from (12) and (21) the dual estimate

$$
\| - a\Delta\ddot{u}_\eta + \ddot{u}_\eta\|_{L_1(I; H^{-2}(\Omega))} \leq c.
\tag{22}
$$

We take the sequence $\{u_k\} \equiv \{u_{\eta_k}\}$, $\eta_k \to 0+$.

After applying the Lemma 1 with the spaces

$$B_0 = L_2(\Omega), \ B = H^{-1}(\Omega), \ B_1 = H^{-2}(\Omega)$$

we obtain the relative compactness of the sequence $\{-a\Delta \dot{u}_k + \dot{u}_k\}$ in $L_2(I; H^{-1}(\Omega))$ and with the help of the test function $\dot{u}_k - \dot{u}$ the crucial strong convergence

$$\dot{u}_k \to \dot{u} \text{ in } L_2(I; \mathring{H}^1(\Omega)). \tag{23}$$

Simultaneously we have the convergences

$$\begin{aligned}
&\dot{u}_k \rightharpoonup \dot{u} \text{ in } L_2(I; \mathring{H}^2(\Omega)), \\
&\dot{u}_k \to \dot{u} \text{ in } L_2(I; W_p^1(\Omega)), \\
&\tfrac{1}{2}\partial_t\Phi(u_k, u_k) + \partial_t L u_k \rightharpoonup \tfrac{1}{2}\partial_t\Phi(u, u) + \partial_t L u \text{ in } L_2(I; W_p^2(\Omega)).
\end{aligned} \tag{24}$$

It can be verified after inserting the test function $z = y - u_k$ in (12) for $y \in \mathcal{K}$, performing the integration by parts in the terms containing $\ddot{u}$, applying the convergences (23), (24), using the definitions of the operators L, Φ in (6), (7) and the weak lower semicontinuity that the limit function u is a solution of the original problem $\mathcal{P}$.

Remark 1. The existence Theorem 1 can be after some modification verified also for another types of boundary conditions.

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